

## 6.4 TIME-DEPENDENT SCHRÖDINGER'S WAVE EQUATION

A plane wave moving in the  $X$ -direction is represented in the exponential form as

$$y = Ae^{i(\omega t - kx)} \quad \dots(i)$$

where  $\omega = 2\pi\nu$  ( $\nu$  = frequency of the wave) and  $k$  the propagation constant =  $\frac{2\pi}{\lambda}$  ( $\lambda$  = wavelength of the wave).

According to quantum theory, the **energy** of a wave

$$E = h\nu$$

and therefore

$$\nu = \frac{E}{h}$$

and

$$\omega t = 2\pi\nu t = 2\pi \frac{E}{h} t = \frac{Et}{h/2\pi} = \frac{Et}{\hbar}$$

and momentum

$$p_x = \frac{h}{\lambda} \quad \text{or} \quad \frac{1}{\lambda} = \frac{p_x}{h}$$

$\therefore$

$$kx = \frac{2\pi}{\lambda} x = 2\pi \frac{p_x x}{h} = \frac{p_x}{h/2\pi} x = \frac{p_x}{\hbar} x$$

where

$$\hbar = \frac{h}{2\pi} = 1.054 \times 10^{-34} \text{ Js.}$$

Substituting in (i) we can express a plane wave moving in the  $X$ -direction as

$$\begin{aligned} y &= Ae^{i\left(\frac{E}{\hbar}t - \frac{p_x x}{\hbar}\right)} \\ &= Ae^{\frac{i}{\hbar}(Et - xp_x)} \quad \dots(ii) \end{aligned}$$

By analogy between photons and electron beams, we assume that the electrons in a beam travelling in the  $X$ -direction with constant momentum  $p$ , and energy  $E$  are controlled by *pilot* waves with a wave function  $\Psi$  which plays the same role for electrons and other particles as the displacement  $y$  for sound waves or electric field vector  $E$  for electro-magnetic waves.

Thus  $\Psi(x, t)$  as a function of  $x$  and  $t$  denotes the amplitude of matter wave at a point  $x$  at a time  $t$  and is given by an equation similar to equation (ii) i.e.,

$$\Psi(x, t) = Ae^{\frac{i}{\hbar}(Et - xp)}$$

However, in wave mechanics, we use the exponential form with a *negative sign* i.e.,

$$\begin{aligned} \Psi(x, t) &= Ae^{-\frac{i}{\hbar}(Et - xp)} \\ &= Ae^{\frac{i}{\hbar}(xp - Et)} \end{aligned}$$

for a wave moving along the positive direction of  $X$ -axis. We represent a wave moving along the negative direction of  $X$ -axis by the equation

$$\Psi(x, t) = Ae^{-\frac{i}{\hbar}(xp - Et)}$$

This merely introduces a phase difference of  $\pi$  from the notation used for sound waves, electromagnetic and other waves.

Differentiating equation (iii) with respect to time and writing only  $\Psi$  for  $\Psi(x, t)$ , we have

$$\frac{\partial \Psi}{\partial t} = -\frac{i}{\hbar} EAe^{\frac{i}{\hbar}(xp - Et)} = -\frac{i}{\hbar} E\Psi \quad \dots(iii)$$

or

$$E\Psi = -\frac{\hbar}{i} \frac{\partial \Psi}{\partial t} \quad \dots(iv)$$

Differentiating equation (iii) with respect to  $x$ , we have

$$\frac{\partial \Psi}{\partial x} = \frac{i}{\hbar} p_x Ae^{\frac{i}{\hbar}(xp - Et)}$$

and

$$\frac{\partial^2 \Psi}{\partial x^2} = -\frac{p_x^2}{\hbar^2} Ae^{\frac{i}{\hbar}(xp - Et)} = -\frac{p_x^2}{\hbar^2} \Psi$$

or

$$p_x^2 \Psi = -\hbar^2 \frac{\partial^2 \Psi}{\partial x^2} \quad \dots(v)$$

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2. Subject - PHYSICS
3. Class - TAC-III
4. Paper - ✓
5. Topic - Time dependent Schrodinger wave equation
6. Date - 19-01-22